

Entropy of a nonuniformly rectilinearly accelerating black hole

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Adopting thin film brick-wall model, we calculate the entropy of a nonuniformly rectilinearly accelerating non-stationary black hole expressed by Kinnersley metric. Because the black hole is accelerated, the event horizon is axisymmetric. The different points of horizon surface may have different temperature. We calculate the temperature and the entropy density at every point of the horizon at first, then we obtain the total entropy through integration, which is proportional to the area of event horizon as the same as the stationary black holes. It is shown that the black hole entropy may be regarded as the entropy of quantum fields just on the surface of event horizon.

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I. INTRODUCTION

Since Bekenstein suggested that the entropy of a black hole is proportional to its surface area, the concerned research work has got much progress [1-3]. The brick-wall model suggested by 't Hooft gives a statistical explanation to the origin of black hole entropy [4,5]. Recently, the brick-wall model is developed to the thin film brick-wall model [6,7]. In this paper, we adopt the thin film brick-wall model to calculate the entropy of a nonuniformly rectilinearly accelerating non-stationary black hole.

Kinnersley has discussed the spacetime of an arbitrarily accelerating point mass [8]. The metric used in this paper is a simple case in which the direction of acceleration doesn't vary, but the magnitude of acceleration and the mass of black hole vary with time. Because the black hole is accelerated, the horizon is axisymmetric. Unlike the spherically symmetric black hole, the different points of horizon surface may have different temperature [9-12]. It is obvious that in this spacetime the thermal equilibrium does not exist in a large region, and the normal brick-wall method encounters difficulty in this situation. The thin film brick-wall model is adopted to overcome this difficulty. Because the spacetime is axisymmetric and the different points of horizon surface may have different temperature, we calculate the entropy density at every point of the horizon at first, then we obtain the total entropy through integration. If we appropriately select relationship between the thickness of film and its distance to event horizon, the thickness vanishes when the distance goes to zero, so the entropy of black hole can be regarded as the entropy of quantum field just on the event horizon.

In Sec. II we give the line element of spacetime and the surface equation of horizon. In Sec. III we use a method which is proposed by Zhao and Dai to study the temperature of the horizon [9,10]. The entropy of black hole is calculated in Sec. IV in detail, and we give conclusion and discussion in Sec. V.

II. METRIC OF NONUNIFORMLY RECTILINEARLY ACCELERATING BLACK HOLE

The metric of an arbitrarily accelerating point mass has been given by Kinnersley [8]

$$ds^2 = [1 - 2ar \cos \theta - r^2(f^2 + h^2 \sin^2 \theta) - 2Mr^{-1}]du^2 + 2dudr + 2r^2 f dud\theta + 2r^2 h \sin^2 \theta dud\varphi - r^2 d\theta^2 - r^2 \sin^2 \theta d\varphi^2, \quad (1)$$

where

$$\begin{aligned} f &= -a(u) \sin \theta + b(u) \sin \varphi + c(u) \cos \varphi, \\ h &= b(u) \cot \theta \cos \varphi - c(u) \cot \theta \sin \varphi. \end{aligned} \quad (2)$$

a, b, c and M are all arbitrary functions of retarded Eddington-Finkelstein coordinate u . a is the magnitude of acceleration. b and c describe the rate of change of the acceleration's direction. In the case of nonuniformly rectilinearly acceleration, we have $b = c = 0$. The metric can be reduced to

$$ds^2 = (1 - 2ar \cos \theta - r^2 f^2 - 2Mr^{-1})du^2 + 2dudr + 2r^2 f dud\theta - r^2 d\theta^2 - r^2 \sin^2 \theta d\varphi^2, \quad (3)$$

where

$$f = -a \sin \theta. \quad (4)$$

Using the advanced Eddington-Finkelstein coordinate v to replace the retarded coordinate u , and adopting signature $(-, +, +, +)$, the metric becomes

$$\begin{aligned} ds^2 = & -(1 - 2ar \cos \theta - r^2 a^2 \sin^2 \theta - 2Mr^{-1})dv^2 + 2dvdr \\ & - 2r^2 a \sin \theta dv d\theta + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi^2. \end{aligned} \quad (5)$$

$\theta = 0$ is the direction of acceleration and the spacetime is axisymmetric. The determinant of metric is

$$g = -r^4 \sin^2 \theta, \quad (6)$$

and the non-zero contravariant components of metric are

$$\begin{aligned} g^{01} &= g^{10} = 1, \\ g^{11} &= 1 - 2ar \cos \theta - 2Mr^{-1}, \\ g^{12} &= g^{21} = a \sin \theta, \\ g^{22} &= r^{-2}, \\ g^{33} &= r^{-2} \sin^2 \theta. \end{aligned} \quad (7)$$

Now let us find the horizon equation of the spacetime represented by Eq. (5). Considering the axial symmetry of spacetime, the surface equation of event horizon can be written as

$$H = H(v, r, \theta) = 0 \quad \text{or} \quad r = r(v, \theta), \quad (8)$$

which should satisfy null surface condition

$$g^{\mu\nu} \frac{\partial H}{\partial x^\mu} \frac{\partial H}{\partial x^\nu} = 0. \quad (9)$$

Substituting Eq. (7) into Eq. (9), we get

$$2 \frac{\partial H}{\partial v} \frac{\partial H}{\partial r} + \left(1 - 2ar \cos \theta - \frac{2M}{r}\right) \left(\frac{\partial H}{\partial r}\right)^2 + 2a \sin \theta \frac{\partial H}{\partial r} \frac{\partial H}{\partial \theta} + \frac{1}{r^2} \left(\frac{\partial H}{\partial \theta}\right)^2 = 0. \quad (10)$$

From Eq. (8) we can get

$$\frac{\partial H}{\partial r} \frac{\partial r}{\partial \theta} + \frac{\partial H}{\partial \theta} = 0, \quad \frac{\partial H}{\partial r} \frac{\partial r}{\partial v} + \frac{\partial H}{\partial v} = 0. \quad (11)$$

Substituting Eq. (11) into Eq. (10), we have

$$-2\dot{r}_H + \left(1 - 2ar_H \cos \theta - \frac{2M}{r_H}\right) - (2a \sin \theta) r'_H + \frac{r'^2_H}{r_H^2} = 0, \quad (12)$$

where

$$\dot{r}_H = \left(\frac{\partial r}{\partial v}\right)_{r=r_H}, \quad r'_H = \left(\frac{\partial r}{\partial \theta}\right)_{r=r_H}. \quad (13)$$

Surface r_H which satisfies Eq. (12) is the event horizon of nonuniformly rectilinearly accelerating black hole.

III. TEMPERATURE OF BLACK HOLE

In this section, we use a method [9-11] which is proposed by Zhao and Dai to study the temperature of the horizon. The calculation of this method is simple and precision. It can be used to calculate temperature of various black holes, including non-spherically symmetric and non-stationary black holes.

This method is based on the Damour-Ruffini's scheme [12]. The essential point is that if we use tortoise coordinate, the radial part of Klein-Gordon equation near the horizon will be reduced to the standard form of wave equation as

$$\frac{\partial^2 \Phi}{\partial r_*^2} + 2 \frac{\partial^2 \Phi}{\partial v \partial r_*} = 0. \quad (14)$$

From Eq.(12) we know that r_H is the function of θ and v on the horizon. So the generalized tortoise coordinate transformation is proposed as [10]

$$\begin{aligned} r_* &= r + \frac{1}{2\kappa} \ln[r - r_H(\theta, v)], \\ v_* &= v - v_0, \\ \theta_* &= \theta - \theta_0. \end{aligned} \quad (15)$$

where κ is an adjustable parameter. Both v_0 and θ_0 are arbitrarily fixed parameter. They are constant under the tortoise coordinate transformation.

From Eq. (15) we get

$$\begin{aligned} \frac{\partial}{\partial r} &= \left[1 + \frac{1}{2\kappa(r - r_H)} \right] \frac{\partial}{\partial r_*}, \\ \frac{\partial}{\partial v} &= \frac{\partial}{\partial v_*} - \frac{\dot{r}_H}{2\kappa(r - r_H)} \frac{\partial}{\partial r_*}, \\ \frac{\partial}{\partial \theta} &= \frac{\partial}{\partial \theta_*} - \frac{r'_H}{2\kappa(r - r_H)} \frac{\partial}{\partial r_*}; \end{aligned} \quad (16)$$

Substituting Eq.(6), Eq.(7) and Eq.(16) into Klein-Gordon equation

$$\frac{1}{\sqrt{-g}} \frac{\partial}{\partial x^\mu} \left(\sqrt{-g} g^{\mu\nu} \frac{\partial \Phi}{\partial x^\nu} \right) - \mu^2 \Phi = 0. \quad (17)$$

When $r \rightarrow r_H(v_0, \theta_0)$, $v \rightarrow v_0$ and $\theta \rightarrow \theta_0$, the equation can be reduced to

$$\alpha \frac{\partial^2 \Phi}{\partial r_*^2} + 2 \frac{\partial^2 \Phi}{\partial r_* \partial v_*} + 2\Omega \frac{\partial^2 \Phi}{\partial r_* \partial \theta_*} - G \frac{\partial \Phi}{\partial r_*} = 0, \quad (18)$$

where

$$\alpha = \lim_{\substack{r \rightarrow r_H(v_0, \theta_0) \\ v \rightarrow v_0, \theta \rightarrow \theta_0}} \frac{\{-2\dot{r}_H + (1 - 2ar \cos \theta - \frac{2M}{r})[2\kappa(r - r_H) + 1] - (2a \sin \theta)r'_H\}r^2[2\kappa(r - r_H) + 1] + r_H'^2}{2\kappa(r - r_H)[2\kappa(r - r_H) + 1]r^2}, \quad (19)$$

$$\Omega = - \left(f + \frac{r'_H}{r_H^2} \right) \Big|_{\substack{v \rightarrow v_0 \\ \theta \rightarrow \theta_0}}, \quad (20)$$

$$G = \left(-\frac{2}{r_H} + 16a \cos \theta + \frac{2M}{r_H^2} - \frac{r_h'^2}{r_H^3} + \frac{r'_H}{r_H^2} \cot \theta \right) \Big|_{\substack{v \rightarrow v_0 \\ \theta \rightarrow \theta_0}}. \quad (21)$$

We select the adjustable parameter κ as

$$\kappa = \frac{1}{2r_H} \frac{\frac{M}{r_H^2} - a \cos \theta - \frac{r_H'^2}{r_H^3}}{\frac{M}{r_H^2} + a \cos \theta + \frac{r_H'^2}{2r_H^3}} \Big|_{\substack{v \rightarrow v_0 \\ \theta \rightarrow \theta_0}}. \quad (22)$$

Then we have $\alpha = 1$, and Eq.(18) can be reduced to

$$\frac{\partial^2 \Phi}{\partial r_*^2} + 2 \frac{\partial^2 \Phi}{\partial r_* \partial v_*} + 2\Omega \frac{\partial^2 \Phi}{\partial r_* \partial \theta_*} - G \frac{\partial \Phi}{\partial r_*} = 0. \quad (23)$$

Separating variables as

$$\Phi = R(r_*)e^{-i\omega v_* + ik_\theta \theta_* + in\phi}, \quad (24)$$

we can verify that the radial wave solutions of Eq.(23) are, respectively,

$$\phi_{\text{in}} = e^{-i\omega v_*}, \quad (25)$$

$$\phi_{\text{out}} = e^{-i\omega v_* + Gr_* + i(2\omega - 2\Omega k_\theta)r_*}. \quad (26)$$

ϕ_{in} is the ingoing wave, while ϕ_{out} is the outgoing wave. Near the event horizon r_H , $r_* = \frac{1}{2\kappa} \ln(r - r_H)$, ϕ_{out} can be rewritten as

$$\psi_{\text{out}} = e^{-i\omega v_*} (r - r_H)^{G/2\kappa} (r - r_H)^{i(\omega - \Omega k_\theta)/\kappa}. \quad (27)$$

It is not analytical at the horizon. By analytical continuation rotating $-\pi$ through the lower-half complex r -plane,

$$(r \rightarrow r_H) \rightarrow |r - r_H|e^{-i\pi} = (r_H - r)e^{-i\pi}, \quad (28)$$

we can extend ϕ_{out} from the outside of the black hole into the inside of the black hole [12]

$$\phi_{\text{out}} \rightarrow \phi'_{\text{out}} = e^{-i\omega v_* + Gr_* + i(2\omega - 2\Omega k_\theta)r_*} e^{-i\pi G/2\kappa} e^{\pi(\omega - \Omega k_\theta)/\kappa} \quad (29)$$

The relative scattering probability of the outgoing wave at the horizon is

$$\left| \frac{\phi_{\text{out}}}{\phi'_{\text{out}}} \right|^2 = e^{-2\pi(\omega - \Omega k_\theta)/\kappa}. \quad (30)$$

Then the spectrum of the Hawking radiation is

$$N_\omega = (e^{\frac{\omega - \Omega k_\theta}{T}} \pm 1)^{-1}, \quad (31)$$

where

$$T = \frac{\kappa}{2\pi} = \frac{1}{2\pi} \cdot \frac{1}{2r_H} \frac{\frac{M}{r_H^2} - a \cos \theta - \frac{r_H'^2}{r_H^3}}{\frac{M}{r_H^2} + a \cos \theta + \frac{r_H'^2}{2r_H^3}}. \quad (32)$$

We can see that the temperature of the nonuniformly rectilinearly accelerating black hole depends not only on the time, but also on the polar angle.

IV. ENTROPY OF BLACK HOLE

In the brick-wall model proposed by 't Hooft [4], the black hole entropy is identified with the entropy of thermal gas of quantum fields excitations outside the event horizon. This needs the thermal equilibrium between the external fields and the black hole. However, this qualification is not satisfied in the case of accelerating black hole, because the black hole is non-stationary and the temperature on the horizon is not uniform as shown in Sec. III.

Recently, a new model, the thin film brick-wall model [6,7] is developed from the original brick-wall model. This model considers that the entropy of a black hole should be only related to quantum gas near its horizon. Due to this opinion and the fact that the density of quantum states near the horizon is divergent, it is natural to take only the quantum field in a thin film near the event horizon into account. If we adopt the thin film brick-wall model, we find that although the global thermal equilibrium is not satisfied, the local thermal equilibrium is always exist. So the difficulty encountered by original brick-wall model is overcome.

Because the spacetime is axisymmetric and the temperature on the horizon is not uniform, we will calculate the entropy density at every point of the horizon at first, then we obtain the total entropy through integration.

In the Sec. II, we have obtained the metric of nonuniformly rectilinearly accelerating black hole shown by Eq.(5) and the horizon equation shown by Eq.(12). Introduce a coordinate transformation [13]

$$R = r - r_H(v, \theta), \quad dR = dr - \dot{r}_H dv - r'_H d\theta. \quad (33)$$

The metric becomes

$$ds^2 = -(-2\dot{r}_H + 1 - 2ar \cos \theta - r^2 a^2 \sin^2 \theta - 2Mr^{-1})dv^2 + 2dv dR \\ - 2(r^2 a \sin \theta - r'_H)dv d\theta + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi^2. \quad (34)$$

The following calculation about entropy is based on this metric. The non-zero contravariant components of the metric are

$$g^{01} = g^{10} = 1, \\ g^{11} = -2\dot{r}_H + 1 - 2ar \cos \theta - 2Mr^{-1} - 2a \sin \theta r'_H + \frac{r'^2_H}{r^2}, \\ g^{12} = g^{21} = a \sin \theta - \frac{r'_H}{r^2}, \\ g^{22} = r^{-2}, \\ g^{33} = r^{-2} \sin^{-2} \theta. \quad (35)$$

$g^{11} = 0$ is just the surface equation of horizon.

The form of metric can be changed to

$$ds^2 = g_{00}dv^2 + 2dv dR + g_{02}dv d\theta + g_{22}d\theta^2 + g_{33}d\varphi^2 \\ = \left(g_{00} - \frac{g_{02}^2}{g_{22}}\right)dv^2 + 2dv dR + g_{22} \left(d\theta + \frac{g_{02}}{g_{22}}dv\right)^2 + g_{33}d\varphi^2 \\ = \hat{g}_{00}dv^2 + 2dv dR + g_{22} \left(d\theta + \frac{g_{02}}{g_{22}}dv\right)^2 + g_{33}d\varphi^2, \quad (36)$$

where

$$\hat{g}_{00} = g_{00} - \frac{g_{02}^2}{g_{22}} = - \left[-2\dot{r}_H + \left(1 - 2ar \cos \theta - \frac{2M}{r}\right) - (2a \sin \theta)r'_H + \frac{r'^2_H}{r^2} \right]. \quad (37)$$

We can get the surface equation of horizon from $\hat{g}_{00} = 0$. If we let

$$\frac{d\theta}{dv} = -\frac{g_{02}}{g_{22}}, \quad (38)$$

the metric can be changed to

$$ds^2 = \hat{g}_{00}dv^2 + 2dv dR + g_{33}d\varphi^2. \quad (39)$$

Like Kerr-Newman Black Hole, we can introduce a dragging angular velocity

$$\Omega = \frac{d\theta}{dv} = -\frac{g_{02}}{g_{22}} = a \sin \theta - \frac{r'_H}{r^2}. \quad (40)$$

Ω is just that shown in Eq.(20) and Eq.(31).

Let us substitute the determinant and the contravariant components of metric described by Eq. (34) into Klein-Gordon equation, which describes the scalar field with mass μ ,

$$\frac{1}{\sqrt{-g}} \frac{\partial}{\partial x^\mu} \left(\sqrt{-g} g^{\mu\nu} \frac{\partial \Phi}{\partial x^\nu} \right) = \mu^2 \Phi. \quad (41)$$

We suppose that the solution has the following form [7,14,15]

$$\Phi = e^{-i(Ev-m\varphi)} G(R, \theta), \quad (42)$$

where

$$G(R, \theta) = e^{iS(R, \theta)}. \quad (43)$$

With WKB approximation, we have

$$g^{11}k_R^2 - 2(E - g^{12}k_\theta)k_R + g^{22}k_\theta^2 + m^2g^{33} + \mu^2 = 0, \quad (44)$$

where

$$k_R = \frac{\partial S}{\partial R}, \quad k_\theta = \frac{\partial S}{\partial \theta}. \quad (45)$$

From Eq. (44) we can obtain the relationship between k_R and k_θ as

$$k_R^+ = \frac{E - g^{12}k_\theta}{g^{11}} + \frac{\sqrt{(E - g^{12}k_\theta)^2 - g^{11}(g^{22}k_\theta^2 + m^2g^{33} + \mu^2)}}{g^{11}}, \quad (46)$$

$$k_R^- = \frac{E - g^{12}k_\theta}{g^{11}} - \frac{\sqrt{(E - g^{12}k_\theta)^2 - g^{11}(g^{22}k_\theta^2 + m^2g^{33} + \mu^2)}}{g^{11}}. \quad (47)$$

Free energy of the system is given by

$$F = - \int dE \frac{\Gamma(E)}{e^{\beta(E - \Omega k_\theta)} - 1}, \quad (48)$$

where $\Gamma(E)$ is the total number of modes whose energy is not greater than E . According to the semiclassical quantization condition and the thin film brick-wall model, we have [4-7,14,15]

$$\begin{aligned} \Gamma(E) &= \frac{1}{4\pi^3} \int dm \int d\theta d\varphi \int dk_\theta \left(\int_\epsilon^{\epsilon+\delta} k_R^+ dR + \int_{\epsilon+\delta}^\epsilon k_R^- dR \right) \\ &= \frac{1}{2\pi^3} \int dm \int d\theta d\varphi \int dk_\theta \int_\epsilon^{\epsilon+\delta} \hat{k}_R dR, \end{aligned} \quad (49)$$

where

$$\begin{aligned} \hat{k}_R &= \frac{\sqrt{(E - g^{12}k_\theta)^2 - g^{11}(g^{22}k_\theta^2 + m^2g^{33} + \mu^2)}}{g^{11}} \\ &= \frac{\sqrt{(E - \Omega k_\theta)^2 - g^{11}(g^{22}k_\theta^2 + m^2g^{33} + \mu^2)}}{g^{11}}. \end{aligned} \quad (50)$$

Introducing $\tilde{E} = E - \Omega k_\theta$, we have

$$\hat{k}_R = \frac{\sqrt{\tilde{E}^2 - g^{11}(g^{22}k_\theta^2 + m^2g^{33} + \mu^2)}}{g^{11}}, \quad (51)$$

$$\Gamma(\tilde{E}) = \frac{1}{2\pi^3} \int dm \int d\theta d\varphi \int dk_\theta \int_\epsilon^{\epsilon+\delta} \hat{k}_R dR, \quad (52)$$

$$F = - \int_0^{+\infty} d\tilde{E} \frac{\Gamma(\tilde{E})}{e^{\beta\tilde{E}} - 1}. \quad (53)$$

The surface density of free energy on horizon can be expressed by

$$\sigma_F = - \int_0^{+\infty} d\tilde{E} \frac{\sigma_\Gamma}{e^{\beta\tilde{E}} - 1}. \quad (54)$$

σ_F and σ_Γ are defined as

$$F = \int \sigma_F dA, \quad \Gamma = \int \sigma_\Gamma dA, \quad (55)$$

where

$$dA = r_H^2 \sin \theta d\theta d\varphi. \quad (56)$$

Now, let us study the integration on k_θ and m in Eq. (52). We use little mass approximation in the process of integration, and the result is

$$\begin{aligned} \Gamma(\tilde{E}) &= \frac{\tilde{E}^3}{6\pi^2} \int d\theta d\varphi \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} (g^{22} g^{33})^{-\frac{1}{2}} dR \\ &= \frac{\tilde{E}^3}{6\pi^2} \int d\theta d\varphi \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} (g_{22} g_{33})^{\frac{1}{2}} dR \\ &\approx \frac{\tilde{E}^3}{6\pi^2} \int dA \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} dR \\ &= \int dA \frac{\tilde{E}^3}{6\pi^2} \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} dR. \end{aligned} \quad (57)$$

Comparing with Eq. (55), we get the surface density of $\Gamma(\tilde{E})$ as

$$\sigma_\Gamma = \frac{\tilde{E}^3}{6\pi^2} \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} dR. \quad (58)$$

The surface density of free energy is given by

$$\begin{aligned} \sigma_F &= - \frac{1}{6\pi^2} \int_0^{+\infty} d\tilde{E} \frac{\tilde{E}^3}{e^{\beta\tilde{E}} - 1} \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} dR \\ &= - \frac{\pi^2}{90\beta^4} \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} dR. \end{aligned} \quad (59)$$

From

$$S = \beta^2 \left. \frac{\partial F}{\partial \beta} \right|_{\beta=\beta_H}, \quad (60)$$

we can obtain the surface density of entropy, as

$$\sigma_S = \frac{4\pi^2}{90\beta_H^3} \int_\epsilon^{\epsilon+\delta} (g^{11})^{-2} dR. \quad (61)$$

Because $g^{11} = 0$ is the surface equation of horizon, g^{11} can be expressed by

$$g^{11} = f(r, \theta)(r - r_H). \quad (62)$$

Substituting Eq. (62) into Eq. (61), we complete the integration on R , as

$$\begin{aligned} \sigma_S &= \frac{4\pi^2}{90\beta_H^3} \int_\epsilon^{\epsilon+\delta} \frac{1}{f^2(r - r_H)^2} dR \\ &\approx \frac{4\pi^2}{90\beta_H^3 f_H^2} \int_\epsilon^{\epsilon+\delta} \frac{1}{R^2} dR \\ &= \frac{4\pi^2}{90\beta_H^3 f_H^2} \frac{\delta}{\epsilon(\epsilon + \delta)}. \end{aligned} \quad (63)$$

From Eq.(12),Eq.(35) and Eq.(62), we can obtain the relationship between f_H and κ , as

$$f_H = 2\kappa(1 - 2\dot{r}_H - 2ar'_H \sin \theta + \frac{2r_H'^2}{r_H^2}). \quad (64)$$

Substituting Eq. (64) into Eq. (63), we get

$$\sigma_S = \frac{4\pi^2}{90\beta_H^3} \frac{1}{[\kappa(1 - 2\dot{r}_H - 2ar'_H \sin \theta + \frac{2r_H'^2}{r_H^2})]^2} \frac{\delta}{\epsilon(\epsilon + \delta)} \frac{1}{4}. \quad (65)$$

Substituting $\beta_H = \frac{2\pi}{\kappa}$ into Eq. (65), we get

$$\sigma_S = \frac{1}{P} \frac{\delta}{\epsilon(\epsilon + \delta)} \frac{1}{4}, \quad (66)$$

where

$$P = 90\beta_H(1 - 2\dot{r}_H - 2ar'_H \sin \theta + \frac{2r_H'^2}{r_H^2})^2. \quad (67)$$

Selecting appropriate cut-off distance ϵ and thickness of thin film δ to satisfy

$$\frac{\delta}{\epsilon(\epsilon + \delta)} = P, \quad (68)$$

we can obtain the surface density of entropy, as

$$\sigma_S = \frac{1}{4}. \quad (69)$$

The total entropy is certainly

$$S = \int \sigma_S dA = \frac{1}{4} A_H. \quad (70)$$

V. CONCLUSIONS

By using the generalized tortoise transformation and the thin film brick-wall model, we have studied the temperature and the entropy of a nonuniformly rectilinearly accelerating black hole. Because the black hole is accelerated, the horizon is axisymmetric and the different points of horizon surface may have different temperature. To overcome the difficulty encountered in normal brick-wall model, we adopt the thin film brick-wall model in which only the local thermal equilibrium is needed. We calculate the entropy density at every point of the horizon at first, then we obtain the total entropy through integration. The result is consistent with the well-known conclusion that the black hole entropy is propotional to its area.

We believe that the black hole entropy should be just identified with the entropy of the quantum field on the 2-dimensional surface of event horizon [16]. The thin film brick-wall model can easily explain this opinion. From Eq. (68) we can get

$$\delta = \frac{P\epsilon^2}{1 - P\epsilon}. \quad (71)$$

That is to say, if the relationship of the thickness δ of thin film and the distance ϵ to horizon satisfies above equation, we can always get Eq. (69). Using Eq. (71) we can get

$$\lim_{\epsilon \rightarrow 0} \frac{\delta}{\epsilon(\epsilon + \delta)} = \lim_{\epsilon \rightarrow 0} \frac{\frac{P\epsilon^2}{1 - P\epsilon}}{\epsilon \left[\epsilon + \frac{P\epsilon^2}{1 - P\epsilon} \right]} = \lim_{\epsilon \rightarrow 0} \frac{\frac{P}{1 - P\epsilon}}{1 + \frac{P\epsilon}{1 - P\epsilon}} = P. \quad (72)$$

This indicates that if appropriate relationship of thickness δ and distance ϵ is selected, when the thin film approaches the horizon, the thickness of the film vanishes and the entropy density expressed by Eq. (69) doesn't change. In this case, the thin film itself becomes horizon, the entropy of thin film is just the entroy of horizon, and the black hole entropy is just identified with the entropy of the quantum field on the event horizon.

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